

## The Dirac Equation Derived using 4-Vectors

It can be shown that spin comes from the Poincaré Symmetry of SR, not from a QM axiom.  
Writing a 4-SpinMomentum then leads naturally to the Dirac Equation.

First, we will do a little pure mathematics:

$$\text{Let } \{ (a^2 - b^2) = (a+b)(a-b) = c^2 \}$$

We can multiply by an arbitrary factor (xy):

$$(a^2 - b^2)(xy) = (a+b)(a-b)(xy) = c^2(xy)$$

If we impose the following extra constraint:

$$(a+b)x = (cy)$$

$$(a-b)y = (cx)$$

Then the separated equations are still true when multiplied together:

$$(a+b)x * (a-b)y = (cy) * (cx) \rightarrow (a+b)(a-b)(xy) = c^2(xy)$$

Now add and subtract the separated equations:

$$(a+b)x + (a-b)y = (cy) + (cx)$$

$$(a+b)x - (a-b)y = (cy) - (cx)$$

Gather terms in {a,b,c}:

$$a(x+y) + b(x-y) = c(x+y)$$

$$a(x-y) + b(x+y) = -c(x-y)$$

Let  $X=(x+y)$  and  $Y=(y-x)=-x-y$ , just a change in variable names

$$aX - bY = cX$$

$$-aY + bX = cY$$

Rearrange:

$$aX - bY = cX$$

$$bX - aY = cY$$

Putting into matrix form:

$$\begin{bmatrix} a - b & [X] \\ [b - a] & [Y] \end{bmatrix} = \begin{bmatrix} c & 0 \\ 0 & c \end{bmatrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix}$$

$$\begin{bmatrix} [a - b] & [X] \\ [b - a] & [Y] \end{bmatrix} = \begin{bmatrix} [c] & [0] \\ [0] & [c] \end{bmatrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix}$$

or

Putting into suggestive matrix form...

$$\begin{pmatrix} [1] & [0] \\ [0] & [-1] \end{pmatrix} \begin{matrix} a + \\ b \end{matrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix} = c \begin{bmatrix} [1] & [0] \\ [0] & [1] \end{bmatrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix}$$

$$\begin{pmatrix} [1] & [0] \\ [0] & [-1] \end{pmatrix} \begin{matrix} a + \\ b \end{matrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix} = c \begin{bmatrix} [1] & [0] \\ [0] & [1] \end{bmatrix} \begin{bmatrix} [X] \\ [Y] \end{bmatrix}$$

And again, to confirm that this matches the original equation:

$$aX - bY = cX$$

$$bX - aY = cY$$

or

$$(a-c)X = (b)Y$$

$$(b)X = (a+c)Y$$

Multiply the terms:

$$(a-c)X(a+c)Y = (b)X(b)Y$$

$$(a-c)(a+c)XY = (b)(b)XY$$

$$(a^2 - c^2)XY = (b^2)XY$$

$$(a^2 - c^2) = b^2$$

$$(a^2 - b^2) = c^2$$

So, mathematically:

$(a^2 - b^2) = c^2$  is the defining equation, which holds equivalently for:

$$\{ (a^2 - b^2)(xy) = (c^2)(xy) \}$$

or

$$\begin{pmatrix} [1 & 0] \\ [0 & -1] \end{pmatrix} a + \begin{pmatrix} [0 & -1] \\ [1 & 0] \end{pmatrix} b \begin{pmatrix} [X] \\ [Y] \end{pmatrix} = c \begin{pmatrix} [1 & 0] \\ [0 & 1] \end{pmatrix} \begin{pmatrix} [X] \\ [Y] \end{pmatrix}$$

It is only the nature of  $\{x,y\}$  and  $\{X,Y\}$  that are different.

[Pauli Matrices](#) are a set of dimensionless 2x2 matrices that may be used in totally classical contexts.

In this regard, they are not a quantum postulate, but one of Poincaré Invariance.

$$\sigma^t = \sigma^0 = \mathbf{I}_2 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad \sigma^x = \sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \sigma^y = \sigma^2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad \sigma^z = \sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

The Pauli Spin Matrices can be written in Tensor notation, with each component itself a dimensionless 2x2 matrix

$$\boldsymbol{\Sigma} = \Sigma^\mu_\nu = \text{Diag}[\sigma^0, \sigma^1, \sigma^2, \sigma^3] = \text{Diag}[\sigma^0, \boldsymbol{\sigma}] \rightarrow \text{Diag}[\sigma^t, \sigma^x, \sigma^y, \sigma^z]$$

In [classical mechanics](#), Pauli matrices are useful in the context of the Cayley-Klein parameters.

The 3D matrix  $\mathbf{P}$  corresponding to the 3-position  $\mathbf{x}$  of a point in space is defined in terms of the above Pauli vector matrix,

$$\mathbf{P} = \mathbf{x} \cdot \boldsymbol{\sigma} = x \sigma^x + y \sigma^y + z \sigma^z$$

The 4D matrix  $\mathbf{P}$  corresponding to the 4-Position  $\mathbf{X}$  of a point in spacetime is defined in terms of the above Pauli vector matrix,

$$\mathbf{P} = \mathbf{X} \cdot \boldsymbol{\Sigma} = ct \sigma^t + x \sigma^x + y \sigma^y + z \sigma^z$$

Suppose now that the vector  $\mathbf{x}=(x_1, x_2, x_3)$  is rotated around an axis with unit vector  $\mathbf{n}=(n_1, n_2, n_3)$  through an angle  $\theta$ .

The transformation matrix  $\mathbf{R}(\theta)$  for rotations about an axis through an angle  $\theta$  may be written in terms of Pauli matrices and the unit matrix. These can be gathered together in tensor notation with the 2D unit matrix  $= \mathbf{I}_2 = \sigma^0$ .

$$\mathbf{R}_x(\theta) = e^{-i\theta x/2} = \sigma^0 \cos(\theta/2) - i \sigma^x \sin(\theta/2) = \begin{bmatrix} \cos(\theta/2) & -i \sin(\theta/2) \\ -i \sin(\theta/2) & \cos(\theta/2) \end{bmatrix}$$

$$\mathbf{R}_y(\theta) = e^{-i\theta y/2} = \sigma^0 \cos(\theta/2) - i \sigma^y \sin(\theta/2) = \begin{bmatrix} \cos(\theta/2) & -\sin(\theta/2) \\ \sin(\theta/2) & \cos(\theta/2) \end{bmatrix}$$

$$\mathbf{R}_z(\theta) = e^{-i\theta z/2} = \sigma^0 \cos(\theta/2) - i \sigma^z \sin(\theta/2) = \begin{bmatrix} \cos(\theta/2) - i \sin(\theta/2) & 0 \\ 0 & \cos(\theta/2) + i \sin(\theta/2) \end{bmatrix} = \begin{bmatrix} e^{-i(\theta/2)} & 0 \\ 0 & e^{+i(\theta/2)} \end{bmatrix}$$

For an angle  $\theta$  about an arbitrary axis  $\mathbf{n}$

$$\mathbf{R}_n(\theta) = e^{-i\theta n/2} = \sigma^0 \cos(\theta/2) - i (\mathbf{n} \cdot \boldsymbol{\sigma}) \sin(\theta/2) = \begin{bmatrix} \cos(\theta/2) - i \sin(\theta/2)n_3 & -\sin(\theta/2) (n_2 + in_1) \\ \sin(\theta/2) (n_2 - in_1) & \cos(\theta/2) + i \sin(\theta/2)n_3 \end{bmatrix}$$

All of the relativistic wave equations can be derived from a common source:  
The relativistic mass-energy relation, including spin, in an Electromagnetic (EM) field.  
Note that this formalism fits well with the Stern-Gerlach experiment.

$$4\text{-Momentum } \mathbf{P} = (E/c, \mathbf{p})$$

$$4\text{-Momentum including Spin } \mathbf{P}_s = \boldsymbol{\Sigma} \cdot \mathbf{P} = \Sigma^{\mu\nu} P^\nu = \eta_{\alpha\beta} \Sigma^{\mu\alpha} P^\beta = P_s^\mu$$

$$\Sigma^{\mu\nu} \text{ is the 4D (1,1)-Pauli Spin-Matrix Tensor} = \text{Diag}[\boldsymbol{\sigma}^0, \boldsymbol{\sigma}]$$

$$\Sigma^{\mu\nu} \text{ is the 4D (2,0)-Pauli Spin-Matrix Tensor} = \text{Diag}[\boldsymbol{\sigma}^0, -\boldsymbol{\sigma}]$$

$$\mathbf{P}_s = \text{Diag}[\boldsymbol{\sigma}^0, -\boldsymbol{\sigma}] \cdot \mathbf{P} = \text{Diag}[\boldsymbol{\sigma}^0, -\boldsymbol{\sigma}] \cdot (E/c, \mathbf{p}) = (\boldsymbol{\sigma}^0 E/c, \boldsymbol{\sigma} \cdot \mathbf{p})$$

$$\mathbf{P}_s = (p_s^0, \mathbf{p}_s) = (\boldsymbol{\sigma}^0 E/c, \boldsymbol{\sigma} \cdot \mathbf{p})$$

with  $\boldsymbol{\sigma}^0$  as an identity matrix  $\mathbf{I}$  of appropriate spin dimension and  $\boldsymbol{\sigma}$  is the Pauli Spin Matrix Vector

$$4\text{-Momentum inc. Spin in External Field } \mathbf{P}_T = (H/c, \mathbf{p}_T) = (E_T/c, \mathbf{p}_T)$$

with:

$$\mathbf{H} = E_T = \text{Hamiltonian} = \text{Total Energy Of System}$$

$$\mathbf{p}_T = \text{Total 3-momentum Of System}$$

$$4\text{-TotalMomentum } \mathbf{P}_T = \mathbf{P} + q\mathbf{A}$$

$$4\text{-Momentum } \mathbf{P} = \mathbf{P}_T - q\mathbf{A}$$

$$4\text{-MomentumIncSpin } \mathbf{P}_s = (p_s^0, \mathbf{p}_s) = (\boldsymbol{\sigma}^0 E/c, \boldsymbol{\sigma} \cdot \mathbf{p}) = (\boldsymbol{\sigma}^0 (E_T/c - q\phi/c), \boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})) : \text{Note each component is a 2x2 matrix}$$

$$\mathbf{P}_s \cdot \mathbf{P}_s = (p_s^0)^2 - (\mathbf{p}_s)^2 = [\boldsymbol{\sigma}^0 (E/c)]^2 - [\boldsymbol{\sigma} \cdot (\mathbf{p})]^2 = [\boldsymbol{\sigma}^0 (E_T/c - q\phi/c)]^2 - [\boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 = (m_0 c)^2 = (E_0/c)^2$$

The 4-TotalMomentum (inc. External Field Minimal-Coupling and Spin)

$$\mathbf{P}_s = \boldsymbol{\Sigma} \cdot \mathbf{P} = \boldsymbol{\Sigma} \cdot (\mathbf{P}_T - q\mathbf{A}) = [\boldsymbol{\sigma}^0 (E_T/c - q\phi/c), \boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]$$

with  $\boldsymbol{\Sigma} = \Sigma^{\mu\nu}$  as the Pauli Spin Matrices, and taking the Einstein summation gives the  $\boldsymbol{\sigma}^0$  and  $\boldsymbol{\sigma}$

$$\mathbf{P}_s \cdot \mathbf{P}_s = (\boldsymbol{\Sigma} \cdot \mathbf{P})^2 = [\boldsymbol{\Sigma} \cdot (\mathbf{P}_T - q\mathbf{A})]^2 = [\boldsymbol{\sigma}^0 (E_T/c - q\phi/c)]^2 - [\boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 = (m_0 c)^2$$

$$(\boldsymbol{\Sigma} \cdot \mathbf{P})^2 = (m_0 c)^2$$

$$(\boldsymbol{\Sigma} \cdot \partial)^2 = -(m_0/\hbar)^2$$

$$(\boldsymbol{\Sigma} \cdot \partial)^2 + (m_0 c/\hbar)^2 = 0$$

$$(\boldsymbol{\Sigma} \cdot (D - (i/\hbar)q\mathbf{A}))^2 + (m_0 c/\hbar)^2 = 0$$

Now, to prove that this "Relativistic Pauli" Energy-Momentum equation can give the Dirac equation

$$\mathbf{P}_s \cdot \mathbf{P}_s = [\boldsymbol{\sigma}^0 (E_T/c - q\phi/c)]^2 - [\boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 = (p_s^0)^2 - (\mathbf{p}_s)^2 = (m_0 c)^2 = (E_0/c)^2$$

$$\mathbf{P}_s \cdot \mathbf{P}_s = \mathbf{I} (E_T/c - q\phi/c)^2 - [\boldsymbol{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 = (p_s^0)^2 - (\mathbf{p}_s)^2 = (m_0 c)^2 = (E_0/c)^2$$

$$\mathbf{P}_s \cdot \mathbf{P}_s = (p_s^0)^2 - (\mathbf{p}_s)^2 = (p_s^0 + \mathbf{p}_s) (p_s^0 - \mathbf{p}_s) = (m_0 c)^2$$

$$\{(p_s^0)^2 - (\mathbf{p}_s)^2\} (xy) = (m_0 c)^2 (xy)$$

From our math proof above, this is equivalent to:

$$\begin{pmatrix} [1 & 0] \\ [0 & -1] \end{pmatrix} p_s^0 + \begin{pmatrix} [0 & -1] \\ [1 & 0] \end{pmatrix} \mathbf{p}_s \begin{pmatrix} ) [X] \\ ) [Y] \end{pmatrix} = (m_0 c) I_2 \begin{pmatrix} [X] \\ [Y] \end{pmatrix}$$

or

$$\begin{pmatrix} [1 & 0] \\ [0 & -1] \end{pmatrix} \boldsymbol{\sigma}^0 p^0 + \begin{pmatrix} [0 & -1] \\ [1 & 0] \end{pmatrix} \boldsymbol{\sigma} \cdot \mathbf{p} \begin{pmatrix} ) [X] \\ ) [Y] \end{pmatrix} = (m_0 c) I_2 \begin{pmatrix} [X] \\ [Y] \end{pmatrix}$$

Putting into highly suggestive matrix form...

$$\begin{pmatrix} [\boldsymbol{\sigma}^0 & 0] \\ [0 & -\boldsymbol{\sigma}^0] \end{pmatrix} p^0 + \begin{pmatrix} [0 & -\boldsymbol{\sigma}] \\ [\boldsymbol{\sigma} & 0] \end{pmatrix} \cdot \mathbf{p} \begin{pmatrix} ) [X] \\ ) [Y] \end{pmatrix} = (m_0 c) I_2 \begin{pmatrix} [X] \\ [Y] \end{pmatrix}$$

let Spinor  $\Psi = \begin{pmatrix} [X] \\ [Y] \end{pmatrix}$  and note that  $\sigma^0 = I_2$

this is equivalent to [Dirac Gamma Matrices](#) (in Dirac Basis)...

$$\begin{pmatrix} [I_2 & 0] \\ [0 & -I_2] \end{pmatrix} p^0 + \begin{pmatrix} [0 & -\boldsymbol{\sigma}] \\ [\boldsymbol{\sigma} & 0] \end{pmatrix} \cdot \mathbf{p} \Psi = (m_0 c) I_2 \Psi$$

$$(\gamma^0 p^0 - \boldsymbol{\gamma} \cdot \mathbf{p}) \Psi = (m_0 c) I \Psi$$

$$(\boldsymbol{\Gamma} \cdot \mathbf{P}) \Psi = (m_0 c) I \Psi$$

$$(\boldsymbol{\Gamma} \cdot \mathbf{P}) = (m_0 c)$$

$$(\Gamma^\mu P_\mu) \Psi = (m_0 c) \Psi$$

$$i\hbar(\Gamma^\mu \partial_\mu) \Psi = (m_0 c) \Psi$$

The [Dirac Relativistic Quantum Equation](#) for spin 1/2 particles

To recap, the SR 4-Vectors have some very simple, fundamental, and Invariant Lorentz Scalar relations between one another:

4-Position	$\mathbf{R} = R^\mu = (ct, \mathbf{r}) \in \langle \text{Event} \rangle \in \langle \text{Time Space} \rangle$	$\rightarrow (E_0/c^2) \leftrightarrow \mathbf{P} = (E/c, \mathbf{p})$
4-Velocity	$\mathbf{U} = U^\mu = \gamma(\mathbf{c}, \mathbf{u}) = (\mathbf{U} \cdot \partial)\mathbf{R} = (d/d\tau)\mathbf{R}$	/
4-Momentum	$\mathbf{P} = P^\mu = (E/c, \mathbf{p}) = (m_0)\mathbf{U}$	$\mathbf{U} = \gamma(\mathbf{c}, \mathbf{u}) \quad (\hbar)\uparrow \downarrow (1/\hbar)$
4-WaveVector	$\mathbf{K} = K^\mu = (\omega/c, \mathbf{k}) = (1/\hbar)\mathbf{P}$	\
4-Gradient	$\partial = \partial^\mu = (\partial_t/c, -\nabla) = (-i)\mathbf{K}$	$\rightarrow (\omega_0/c^2) \leftrightarrow \mathbf{K} = (\omega/c, \mathbf{k})$

Analysis of Dirac's Constant ( $\hbar = h/2\pi$ ): Planck's Constant ( $h = 2\pi\hbar$ ), and  $(E_0/\omega_0)$  in the context of SRQM:

It is an empirical (observational) fact that the Lorentz Scalar Invariant  $(E_0/\omega_0) = (\gamma E_0/\gamma \omega_0) = (E/\omega) \implies (\hbar) \quad [J \cdot s] / \{\text{rad}\}$  for all known experimental measurements. The SR 4D-Tensor rules show that one doesn't need a quantum axiom for this.

$(\hbar)$  is actually an empirically-measurable quantity, just like  $(c)$ ,  $(e)$ ,  $(G)$ ,  $(k_B)$ ,  $(\mu_0)$ ,  $(\epsilon_0)$  or the other fundamental constants, which are also 4D Lorentz Scalar Invariants.  $(\hbar)$  can be measured classically {without need of quantum axioms} from the photoelectric effect, from the inverse photoelectric effect, from LED's (injection electroluminescence), from Atomic Line Spectra (Bohr Atom), from the Duane-Hunt Law in Bremsstrahlung, from Electron Diffraction in crystals, from the Watt/Kibble-Balance, from the Sagnac Effect, from Incandescent Blackbody Intensity-Temperature Relations, from Compton Scattering, from the Stern-Gerlach experiment, etc.

For physics simulations of measurements of Dirac's: Planck's Constant ( $\hbar : h$ ), see:

<a href="http://scirealm.org/Physics-PlanckConstantViaAtomicLineSpectra.html">http://scirealm.org/Physics-PlanckConstantViaAtomicLineSpectra.html</a>	(transitions of electron energy-levels)
<a href="http://scirealm.org/Physics-PlanckConstantViaComptonScattering.html">http://scirealm.org/Physics-PlanckConstantViaComptonScattering.html</a>	(photon:electron collisions)
<a href="http://scirealm.org/Physics-PlanckConstantViaElectronDiffraction.html">http://scirealm.org/Physics-PlanckConstantViaElectronDiffraction.html</a>	(electron:atomic crystal scattering)
<a href="http://scirealm.org/Physics-PlanckConstantViaGravityInducedInterferometry.html">http://scirealm.org/Physics-PlanckConstantViaGravityInducedInterferometry.html</a>	(gravity interaction)
<a href="http://scirealm.org/Physics-PlanckConstantViaIncandescence.html">http://scirealm.org/Physics-PlanckConstantViaIncandescence.html</a>	(temperature interaction)
<a href="http://scirealm.org/Physics-PlanckConstantViaLEDs.html">http://scirealm.org/Physics-PlanckConstantViaLEDs.html</a>	(electron:photon interaction)
<a href="http://scirealm.org/Physics-PlanckConstantViaSagnacEffect.html">http://scirealm.org/Physics-PlanckConstantViaSagnacEffect.html</a>	(rotation interaction)
<a href="http://scirealm.org/Physics-PlanckConstantViaSternGerlach.html">http://scirealm.org/Physics-PlanckConstantViaSternGerlach.html</a>	(magnetic-field interaction)

$\{E_0/\omega_0 = E/\omega = \hbar\}$  implies the following SR ( wave:particle-like  $\square \cdot \cdot$ ) relation:

$$\mathbf{P} = \hbar\mathbf{K} = (E/c, \mathbf{p}) = \hbar(\omega/c, \mathbf{k}) = (\hbar\omega/c, \hbar\mathbf{k})$$

The **temporal** part  $\{E = \hbar\omega\}$  gives Einstein's photoelectric quantum "postulate". Again, emphasis: derived from SR.

The **spatial** part  $\{\mathbf{p} = \hbar\mathbf{k}\}$  gives de Broglie's matter-wave quantum "postulate". Again, emphasis: derived from SR.

$$\partial = -i\mathbf{K} = (\partial_t/c, -\nabla) = -i(\omega/c, \mathbf{k})$$

The **temporal** part  $\{\partial_t = -i\omega\}$  or  $\{\omega = i\partial_t\}$  gives temporal:frequency complex planewave change:operator

The **spatial** part  $\{\nabla = i\mathbf{k}\}$  or  $\{\mathbf{k} = -i\nabla\}$  gives spatial:wavenumber complex planewave change:operator

The standard Schrödinger QM Relations derived from SR:

Again, we examine these SR 4-Vector relations derived above... and by simply combining them...

$$\mathbf{P} = \hbar\mathbf{K} \quad (\text{a relation which is entirely empirical, based on just SR arguments, shown above})$$

$$\mathbf{K} = i\partial \quad (\text{which is a relation for complex plane-waves, used in classical EM and elsewhere})$$

$$\mathbf{P} = \hbar i\partial = (E/c, \mathbf{p}) = \hbar i(\partial_t/c, -\nabla)$$

The **temporal** part  $\{E = \hbar i\partial_t = \hbar \partial/\partial t\}$  gives unitary QM time evolution operator. Again, emphasis: derived from SR.

The **spatial** part  $\{\mathbf{p} = -\hbar i\nabla\}$  gives the QM momentum operator. Again, emphasis: derived from SR.

Relativistic version (Relativistic Pauli Equation)

$$\begin{aligned} \mathbf{P}_s \cdot \mathbf{P}_s &= [\underline{\sigma}^0 (E_T/c - q\phi/c)]^2 - [\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 = (m_0c)^2 = (E_0/c)^2 \\ [\underline{\sigma}^0 (E_T/c - q\phi/c)]^2 - [\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\mathbf{I}} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - i[\underline{\sigma} \cdot ((\mathbf{p}_T \times -q\mathbf{a}) + (-q\mathbf{a} \times \mathbf{p}_T))] &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\mathbf{I}} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - i[\underline{\sigma} \cdot i\hbar q((-\nabla_T \times -\mathbf{a}) + (-\mathbf{a} \times -\nabla_T))] &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\mathbf{I}} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - i[\underline{\sigma} \cdot i\hbar q(\mathbf{B})] &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\mathbf{I}} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 + \hbar q[\underline{\sigma} \cdot \mathbf{B}] &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 = [\underline{\mathbf{I}} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - \hbar q[\underline{\sigma} \cdot \mathbf{B}] + (m_0c)^2 \end{aligned}$$

Non-relativistic version (Standard [Pauli Equation](#))

$$\begin{aligned} [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 - [\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 &= (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)]^2 = [\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 + (m_0c)^2 \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)] &= (\pm) \text{Sqrt} [[\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 + (m_0c)^2] \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)] &= (\pm) \text{Sqrt} [[\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 + (m_0c)^2] \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)] &\sim (\pm) [[\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 / (2m_0c) + (m_0c)] \\ [\underline{\mathbf{I}}(E_T/c - q\phi/c)] &\sim (\pm) [([\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - \hbar q[\underline{\sigma} \cdot \mathbf{B}]) / (2m_0c) + (m_0c)] \\ [\underline{\mathbf{I}}(E_T - q\phi)] &\sim (\pm) [([\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - \hbar q[\underline{\sigma} \cdot \mathbf{B}]) / (2m_0) + (m_0c^2)] \\ E_T \sim q\phi &(\pm) [([\underline{\sigma} \cdot (\mathbf{p}_T - q\mathbf{a})]^2 - \hbar q[\underline{\sigma} \cdot \mathbf{B}]) / (2m_0) + (m_0c^2)] \\ \text{where the } -(\hbar q[\underline{\sigma} \cdot \mathbf{B}]) / (2m_0) &\text{ is the [Stern-Gerlach term](#)} \end{aligned}$$